

Describing M5-branes using Loop Spaces

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Based on:

- CS, [arXiv:1007.3301](#), CMP ...
- C. Papageorgakis and CS, [arXiv:1103.6192](#), JHEP ...
- Sam Palmer and CS [arXiv:1105.????](#)

Motivation

Find an algorithm for the construction of self-dual string solutions

- **Effective description of M2-branes** proposed in 2007.
- This created lots of interest:
BLG-model: >467 citations, **ABJM-model**: >605 citations
- Inspired by an idea by **Basu-Harvey**:
Propose a lift of the **Nahm eqn.** describing D1-D3-system:
Basu-Harvey eqn. describes M2-M5-brane system
- **Nahm** transform:
go from Nahm eqn. to **Bogomolny monopole eqn.**
switch perspective from D1-brane to D3-brane
- **Is there a lift for this Nahm transform?**
go from BH eqn. to self-dual string eqn.
switch perspective from M2-brane to M5-brane
- Such a transform would open up interesting possibilities:
eff. description of **M5-branes**, new **integrable** structures, ...
- Alternative Approach: **Gustavsson, 0802.3456**

Outline

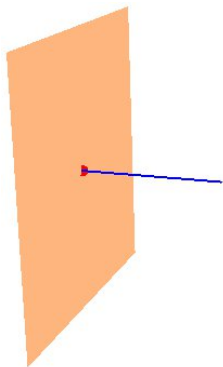
We will discuss the construction of monopoles and lift each ingredient to M-theory.

- Basu-Harvey lift of the Nahm equation and 3-Lie algebras
- Monopoles and self-dual strings
- Principal $U(1)$ -bundles, abelian gerbes and loop space
- ADHMN construction and its lift
- Examples of self-dual string solutions
- Non-abelian tensor multiplet on loop space

D1-D3-Branes and the Nahm Equation

D1-branes ending on D3-branes can be described by the Nahm equation.

dim	0	1	2	3	...	6
D1	×					×
D3	×	×	×	×		



N D1-branes ending on D3-branes:

A **Monopole** appears.

$X^i \in \mathfrak{u}(N)$: transverse fluctuations

Nahm equation: ($s = x^6$)

$$\frac{d}{ds} X^i + \varepsilon^{ijk} [X^j, X^k] = 0$$

Note $SO(3)$ -invariance.

Solution: $X^i = r(s)G^i$ with

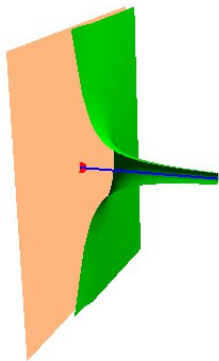
$$r(s) = \frac{1}{s}, \quad G^i = \varepsilon^{ijk} [G^j, G^k]$$

Nahm, Diaconescu, Tsimpis

D1-D3-Branes and the Nahm Equation

The D1-branes end on the D3-branes by forming a fuzzy funnel.

dim	0	1	2	3	...	6
D1	×					×
D3	×	×	×	×		



Solution: $X^i = r(s)G^i$

$$r(s) = \frac{1}{s}, \quad G^i = \varepsilon^{ijk} [G^j, G^k]$$

The D1-branes form a **fuzzy funnel**:

G^i form irrep of $\mathfrak{su}(2)$:

coordinates on fuzzy sphere S_F^2

D1-worldvolume polarizes: $2d \rightarrow 4d$

Myers

Lifting D1-D3-Branes to M2-M5-Branes

The lift to M-theory is performed by a T-duality and an M-theory lift

IIB	0	1	2	3	4	5	6
D1	×						×
D3	×	×	×	×			

T-dualize along x^5 :

IIA	0	1	2	3	4	5	6
D2	×					×	×
D4	×	×	×	×		×	

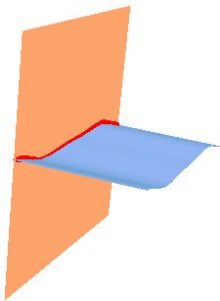
Interpret x^4 as M-theory direction:

M	0	1	2	3	4	5	6
M2	×					×	×
M5	×	×	×	×	×	×	

The Basu-Harvey lift of the Nahm Equation

M2-branes ending on M5-branes yield a Nahm equation with a cubic term.

M	0	1	2	3	4	5	6
M2	×					×	×
M5	×	×	×	×	×	×	



Basu, Harvey, hep-th/0412310

A **Self-Dual String** appears.

Substitute **SO(3)**-inv. **Nahm eqn.**

$$\frac{d}{ds} X^i + \varepsilon^{ijk} [X^j, X^k] = 0$$

by the **SO(4)**-invariant equation

$$\frac{d}{ds} X^\mu + \varepsilon^{\mu\nu\rho\sigma} [X^\nu, X^\rho, X^\sigma] = 0$$

Solution: $X^\mu = r(s)G^\mu$ with

$$r(s) = \frac{1}{\sqrt{s}}, \quad G^\mu = \varepsilon^{\mu\nu\rho\sigma} [G^\nu, G^\rho, G^\sigma]$$

The Basu-Harvey lift of the Nahm Equation

M2-branes ending on M5-branes yield a Nahm equation with a cubic term.

M	0	1	2	3	4	5	6
M2	×					×	×
M5	×	×	×	×	×	×	

Solution: $X^\mu = r(s)G^\mu$

$$r(s) = \frac{1}{\sqrt{s}}, \quad G^\mu = \varepsilon^{\mu\nu\rho\sigma} [G^\nu, G^\rho, G^\sigma]$$

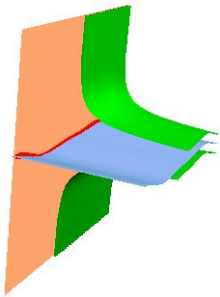
The M2-branes form a **fuzzy funnel**:

G^μ form a rep of $\mathfrak{so}(4)$:

coordinates on fuzzy sphere S_F^3

M2-worldvolume polarizes: $3d \rightarrow 6d$

What is this triple bracket?



What is the algebra behind the triple bracket?

In analogy with Lie algebras, we can introduce 3-Lie algebras.

$$\frac{d}{ds} X^\mu + [A_s, X^\mu] + \varepsilon^{\mu\nu\rho\sigma} [X^\nu, X^\rho, X^\sigma] = 0, \quad X^\mu \in \mathcal{A}$$

Trivial: \mathcal{A} is a **vector space**, $[\cdot, \cdot, \cdot]$ **trilinear+antisymmetric**.

▷ Gauge transformations from **inner derivations**:

Triple bracket forms a map $D : \mathcal{A} \wedge \mathcal{A} \rightarrow \text{Der}(\mathcal{A}) =: \mathfrak{g}_{\mathcal{A}}$ via

$$D(A, B) \triangleright C := [A, B, C]$$

Demand a “3-Jacobi identity,” the **fundamental identity**:

$$\begin{aligned} D(A, B) \triangleright (D(C, D) \triangleright E) &:= [A, B, [C, D, E]] \\ &= [[A, B, C], D, E] + [C, [A, B, D], E] + [C, D, [A, B, E]] \end{aligned}$$

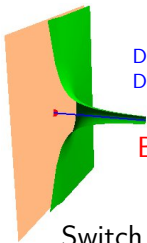
The inner derivations form indeed a **Lie algebra**:

$$[D(A, B), D(C, D)] \triangleright E := D(A, B) \triangleright (D(C, D) \triangleright E) - D(C, D) \triangleright (D(A, B) \triangleright E)$$

Bracket closes due to **fundamental identity**.

Monopoles and Self-Dual Strings

Lifting monopoles to M-theory yields self-dual strings.



	0	1	2	3	4	5	6
D1	×						×
D3	×	×	×	×			

BPS configuration!

Switch perspective: $D1 \rightarrow D3$:

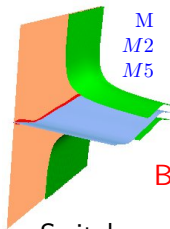
Bogomolny monopole eqn.:

$$F_{ij} = \varepsilon_{ijk} \nabla_k \Phi \Rightarrow \nabla^2 \Phi = 0$$

Single D3: **Dirac monopole**

$$\Phi = \frac{1}{r} \Rightarrow r(s) = \frac{1}{s}$$

\Rightarrow **matching profile!**



M	0	1	2	3	4	5	6
M2	×					×	×
M5	×	×	×	×	×	×	

BPS configuration!

Switch perspective: $M2 \rightarrow M5$:

Self-dual string eqn.:

$$H_{\mu\nu\rho} = \varepsilon_{\mu\nu\rho\sigma} \partial_\sigma \Phi \Rightarrow \partial^2 \Phi = 0$$

Only single M5 known:

$$\Phi = \frac{1}{r^2} \Rightarrow r(s) = \frac{1}{\sqrt{s}}$$

\Rightarrow **matching profile!**

Dirac Monopoles and Principal $U(1)$ -bundles

Dirac monopoles are described by principal $U(1)$ -bundles over S^2 .

Manifold M with cover $(U_i)_i$. **Principal $U(1)$ -bundle** over M :

$$F \in \Omega^2(M, \mathfrak{u}(1)) ,$$

$$A_{(i)} \in \Omega^1(U_i, \mathfrak{u}(1)) \text{ with } F = dA_{(i)}$$

$$g_{ij} \in \Omega^0(U_i \cap U_j, U(1)) \text{ with } A_{(i)} - A_{(j)} = d \log g_{ij}$$

Consider monopole in \mathbb{R}^3 , **but** describe it on S^2 around monopole:

S^2 with patches U_+, U_- , $U_+ \cap U_- \sim S^1$: $g_{+-} = e^{-in\phi}$, $n \in \mathbb{Z}$

$$c_1 = \frac{i}{2\pi} \int_{S^2} F = \frac{i}{2\pi} \int_{S^1} A^+ - A^- = \frac{1}{2\pi} \int_0^{2\pi} n d\phi = n$$

Monopole charge: n

Self-Dual Strings and Abelian Gerbes

Self-dual strings are described by abelian gerbes.

Manifold M with cover $(U_i)_i$. **Abelian (local) gerbe** over M :

$$H \in \Omega^3(M, \mathfrak{u}(1)) ,$$

$$B_{(i)} \in \Omega^2(U_i, \mathfrak{u}(1)) \text{ with } H = dB_{(i)}$$

$$A_{(ij)} \in \Omega^1(U_i \cap U_j, \mathfrak{u}(1)) \text{ with } B_{(i)} - B_{(j)} = dA_{ij}$$

$$h_{ijk} \in \Omega^0(U_i \cap U_j \cap U_k, \mathfrak{u}(1)) \text{ with } A_{(ij)} - A_{(ik)} + A_{(jk)} = dh_{ijk}$$

Note: Local gerbe: principal $U(1)$ -bundles on intersections $U_i \cap U_j$.

Consider S^3 , patches $U_+, U_-, U_+ \cap U_- \sim S^2$: **bundle over S^2**

Reflected in: $H^2(S^2, \mathbb{Z}) \cong H^3(S^3, \mathbb{Z}) \cong \mathbb{Z}$

$$\frac{i}{2\pi} \int_{S^3} H = \frac{i}{2\pi} \int_{S^2} B_+ - B_- = \dots = n$$

Charge of self-dual string: n

Abelian Gerbes and loop space

By going to loop space, one can reduce differential forms by one degree.

Consider the following **double fibration**:

$$\begin{array}{ccc} & \mathcal{L}M \times S^1 & \\ ev \swarrow & & \searrow pr \\ M & & \mathcal{L}M \end{array}$$

Identify $T\mathcal{L}M = \mathcal{L}TM$, then: $x \in \mathcal{L}M \Rightarrow \dot{x}(\tau) \in T\mathcal{L}M$

Transgression

$$\mathcal{T} : \Omega^{k+1}(M) \rightarrow \Omega^k(\mathcal{L}M) , \quad \mathcal{T} = pr^! \circ ev^*$$

$$(\mathcal{T}\omega)_x(v_1(\tau), \dots, v_k(\tau)) := \int_{S^1} d\tau \omega(v_1(\tau), \dots, v_k(\tau), \dot{x}(\tau))$$

An abelian local gerbe over M is a principal $U(1)$ -bundle over $\mathcal{L}M$.

Note: We will work on $\mathcal{L}M \times S^1$, can do $\mathcal{L}M$, too.

The ADHMN construction

There is a map between and solutions to the Nahm equations.

Nahm transform: Instantons on $T^4 \mapsto$ instantons on $(T^4)^*$

Roughly here:

$$T^4: \begin{cases} 3 \text{ rad. } 0 \\ 1 \text{ rad. } \infty : \text{ D1 WV} \end{cases} \quad \text{and} \quad (T^4)^*: \begin{cases} 3 \text{ rad. } \infty : \text{ D3 WV} \\ 1 \text{ rad. } 0 \end{cases}$$

Introduce (twisted) “**Dirac operators**”:

$$\nabla_{s,x} = -\mathbb{1} \frac{d}{ds} + \sigma^i \otimes (iX^i + x^i \mathbb{1}_k), \quad \bar{\nabla}_{s,x} := \mathbb{1} \frac{d}{ds} + \sigma^i \otimes (iX^i + x^i \mathbb{1}_k)$$

Properties:

$$\Delta_{s,x} := \bar{\nabla}_{s,x} \nabla_{s,x} > 0, \quad [\Delta_{s,x}, \sigma^i] = 0 \Leftrightarrow X^i \text{ satisfy Nahm eqn.}$$

Normalized **zero modes**: $\bar{\nabla}_{s,x} \psi_{s,x,\alpha} = 0$, $\mathbb{1} = \int_{\mathcal{I}} ds \bar{\psi}_{s,x} \psi_{s,x}$ yield:

$$A_\mu := \int_{\mathcal{I}} ds \bar{\psi}_{s,x} \frac{\partial}{\partial x^\mu} \psi_{s,x} \quad \text{and} \quad \Phi := -i \int_{\mathcal{I}} ds \bar{\psi}_{s,x} s \psi_{s,x}$$

This is a solution to the Bogomolny monopole equations!

Examples: Dirac monopoles

One can easily construct Dirac monopole solutions using the ADHMN construction.

Charge 1: Nahm eqn: $\partial_s X^i = 0$, so put $X^i = 0$. Zero mode:

$$\psi_+ = e^{-sR} \frac{\sqrt{R+x^3}}{x^1 - ix^2} \begin{pmatrix} x^1 - ix^2 \\ R - x^3 \end{pmatrix}$$

Monopole solution:

$$\Phi^+ = -\frac{i}{2R}, \quad A_i^+ = \frac{i}{2(x^1+x^2)^2} \left(x^2 \left(1 - \frac{x^3}{R} \right), -x^1 \left(1 - \frac{x^3}{R} \right), 0 \right)$$

Charge 2: Nahm eqn. nontrivial. Choose:

$$X^i = -\frac{1}{s} T^i \quad \text{with} \quad T^i = \frac{\sigma^i}{2i} = -\bar{T}^i$$

Resulting solution:

$$\Phi^+ = -\frac{i}{R}, \quad A_i^+ = \dots$$

Lift of the “Dirac operator”

There is a natural lift of the Dirac operator to M-theory.

Type IIB (twisted):

$$\nabla_{s,x}^{\text{IIB}} = -\mathbb{1} \frac{d}{ds} + \sigma^i (iX^i + x^i \mathbb{1}_k)$$

IIB	0	1	2	3	4	5	6
<i>D1</i>	×						×
<i>D3</i>	×	×	×	×			

Type IIA (twisted):

$$\nabla_{s,x}^{\text{IIA}} = -\gamma_5 \mathbb{1}_k \frac{d}{ds} + \gamma^4 \gamma^i (X^i - ix^i)$$

IIA	0	1	2	3	4	5	6
<i>D2</i>	×					×	×
<i>D4</i>	×	×	×	×		×	

M-theory (untwisted):

$$\nabla_s^{\text{M}} = -\gamma_5 \frac{d}{ds} + \frac{1}{2} \gamma^{\mu\nu} D(X^\mu, X^\nu)$$

M	0	1	2	3	4	5	6
<i>M2</i>	×					×	×
<i>M5</i>	×	×	×	×	×	×	

M-theory (twisted):

$$\nabla_{s,x(\tau)}^{\text{M}} = -\gamma_5 \frac{d}{ds} + \gamma^{\mu\nu} \left(\frac{1}{2} D(X^\mu, X^\nu) - ix^\mu(\tau) \dot{x}^\nu(\tau) \right)$$

Lifted ADHMN Construction

The lifted ADHMN construction yields solutions to the loop space self-dual string eqns.

Recall: $\Delta^{\text{IIB}} := \bar{\nabla}^{\text{IIB}} \nabla^{\text{IIB}}$, $[\Delta^{\text{IIB}}, \sigma^i] = 0 \Leftrightarrow X^i$ satisfy Nahm eqn.

Here: $\Delta^{\text{M}} := \bar{\nabla}^{\text{M}} \nabla^{\text{M}}$, $[\Delta, \gamma^{\mu\nu}] = 0 \Leftrightarrow X^\mu$ satisfy BH eqn.

Our Dirac operator involved **loop space**, so we need to **transgress**:

$$H = \left(\varepsilon_{\mu\nu\rho\sigma} \frac{\partial}{\partial x^\sigma} \Phi \right) dx^\mu \wedge dx^\nu \wedge dx^\rho$$

is turned into

$$F_{\mu\nu}(x(\tau)) := \frac{\partial}{\partial x^{[\mu}} A_{\nu]}(x(\tau)) = \varepsilon_{\mu\nu\rho\sigma} \dot{x}^\rho(\tau) \frac{\partial}{\partial x^\sigma} \Phi(x(\tau))$$

From normalized, **A-valued** zero modes $\psi_{s,x(\tau)}$ of $\bar{\nabla}^{\text{M}}$ construct

$$A_\mu = \int ds \bar{\psi}_{s,x(\tau)} \frac{\partial}{\partial x^\mu} \psi_{s,x(\tau)}, \quad \Phi = -i \int ds \bar{\psi}_{s,x(\tau)} s \psi_{s,x(\tau)}$$

Verification of the Construction

Verifying the construction is rather straightforward.

$$\begin{aligned}
 F_{\mu\nu} &= \int ds (\partial_{[\mu} \bar{\psi}_s) \partial_{\nu]} \psi_s \\
 &= \int ds \int dt (\partial_{[\mu} \bar{\psi}_s) (\psi_s \bar{\psi}_t - \nabla_s^M G^M(s, t) \bar{\nabla}_t^M) \partial_{\nu]} \psi_t \\
 &= \int ds \int dt \bar{\psi}_s \left(\gamma^{\mu\kappa} \dot{x}^\kappa G^M(s, t) \gamma^{\nu\lambda} \dot{x}^\lambda - \gamma^{\nu\kappa} \dot{x}^\kappa G^M(s, t) \gamma^{\mu\lambda} \dot{x}^\lambda \right) \psi_t
 \end{aligned}$$

Identity : $[\gamma^{\mu\kappa}, \gamma^{\nu\lambda}] \dot{x}^\kappa \dot{x}^\lambda = -2\varepsilon_{\mu\nu\rho\sigma} \gamma^{\sigma\kappa} \gamma_5 \dot{x}^\rho \dot{x}^\kappa$

$$\begin{aligned}
 F_{\mu\nu} &= -\varepsilon_{\mu\nu\rho\sigma} \int ds \int dt \bar{\psi}_s (2\gamma^{\sigma\kappa} \gamma_5 G^M(s, t) \dot{x}^\rho \dot{x}^\kappa) \psi_t \\
 &= -i\varepsilon_{\mu\nu\rho\sigma} \dot{x}^\rho \int ds \int dt \left((\partial_\sigma \bar{\psi}_s) (\psi_s \bar{\psi}_t - \nabla_s^M G^M(s, t) \bar{\nabla}_t^M) t \psi_t + \right. \\
 &\quad \left. \bar{\psi}_s s (\psi_s \bar{\psi}_t - \nabla_s^M G^M(s, t) \bar{\nabla}_t^M) \partial_\sigma \psi_t \right) \\
 &= -i\varepsilon_{\mu\nu\rho\sigma} \dot{x}^\rho \int ds (\partial_\sigma \bar{\psi}_s) s \psi_s + \bar{\psi}_s s \partial_\sigma \psi_s \\
 &= \varepsilon_{\mu\nu\rho\sigma} \dot{x}^\rho \partial_\sigma \Phi
 \end{aligned}$$

Reduction to the ADHMN Construction

The lift reduces in the expected way to the ADHMN construction.

On $\mathcal{L}S^3 \subset \mathcal{L}\mathbb{R}^4$: $x^\mu x^\mu = \dot{x}^\mu \dot{x}^\mu = R^2$, $x^\mu \dot{x}^\mu = 0$.

Reduction (cf. Mukhi/Papageorgakis, 0803.3218):

$$\langle X^4 \rangle = \frac{r}{\ell_p^{3/2}} e_4 = g_{\text{YM}} e_4, \quad \dot{x}^4(\tau_0) = R \Rightarrow \dot{x}^i(\tau_0) = x^4(\tau_0) = 0$$

$$F_{\mu\nu} = \varepsilon_{\mu\nu\rho\sigma} \dot{x}^\rho \frac{\partial}{\partial x^\sigma} \Phi_{\text{SDS}} \quad \rightarrow \quad F_{ij} = \varepsilon_{ijk} \frac{\partial}{\partial x^k} R \Phi_{\text{SDS}} + \dots$$

$$\frac{d}{ds} X^\mu = \frac{1}{3!} \varepsilon^{\mu\nu\rho\sigma} [X^\nu, X^\rho, X^\sigma] \quad \rightarrow \quad \frac{d}{ds} X^i = \frac{1}{2} \varepsilon^{ijk} R [X^j, X^k] + \dots$$

$$\begin{aligned} \nabla^{\text{M}} &= -\gamma_5 \frac{d}{ds} + \gamma^{\mu\nu} \left(\frac{1}{2} D(X^\mu, X^\nu) - i x^\mu(\tau) \dot{x}^\nu(\tau) \right) \\ &\rightarrow -\gamma_5 \frac{d}{ds} + \gamma^{\mu\nu} \left(\frac{1}{2} D(X^\mu, X^\nu) - i x^\mu(\tau_0) \dot{x}^\nu(\tau_0) \right) \\ &= -\gamma_5 \frac{d}{ds} + R \gamma^{4i} \left(X^{i\alpha} D(e_\alpha, e_4) - i x^i(\tau_0) \right) + \dots = \nabla^{\text{IIA}} + \dots \end{aligned}$$

Examples

Our examples reproduce the expected solutions.

Charge 1: Choose again **trivial Nahm data**. Zero modes:

$$\psi \sim e^{-R^2 s} \begin{pmatrix} i(R^2 + x^2 \dot{x}^1 - x^1 \dot{x}^2 - x^4 \dot{x}^3 + x^3 \dot{x}^4) \\ x^3(\dot{x}^1 + i\dot{x}^2) + x^4(\dot{x}^2 - i\dot{x}^1) - (x^1 + ix^2)(\dot{x}^3 - i\dot{x}^4) \\ 0 \\ 0 \end{pmatrix}$$

Solution:

$$\Phi = \frac{i}{2R^2}, \quad F = \frac{2i \sin \theta^1 \sin^2 \theta^2 (\dot{\theta}^2 d\phi \wedge d\theta^1 - \dot{\theta}^1 d\phi \wedge d\theta^2 + \dot{\phi} d\theta^1 \wedge d\theta^2)}{\sqrt{\dot{\phi}^2 + 2(\dot{\theta}^1)^2 + 4(\dot{\theta}^2)^2 - (\dot{\phi}^2 + 2(\dot{\theta}^1)^2) \cos(2\theta^2) - 2\dot{\phi}^2 \cos(2\theta^1) \sin^2 \theta^2}}$$

This solves the loop-space self-dual string equation.

Regression:

$$\begin{aligned} H &= F|_{\dot{\theta}^1=1, \dot{\theta}^2=0, \dot{\phi}=0} \wedge \sin \theta^2 d\theta^1 - F|_{\dot{\theta}^1=0, \dot{\theta}^2=1, \dot{\phi}=0} \wedge d\theta^2 \\ &\quad + F|_{\dot{\theta}^1=0, \dot{\theta}^2=0, \dot{\phi}=1} \wedge \sin \theta^1 \sin \theta^2 d\phi \\ &= 6i \sin \theta^1 \sin^2 \theta^2 d\theta^1 \wedge d\theta^2 \wedge d\phi, \end{aligned}$$

This is indeed the expected solution.

Examples

Our examples reproduce the expected solutions.

Charge 2:

Nahm data:

$$X^\mu = \frac{e_\mu}{\sqrt{2s}}, \quad e_\mu \text{ generate } \mathcal{A}$$

Solution:

$$\Phi(x) = \frac{i}{R^2}$$

As expected: twice the charge of the case $k = 1$.

Remarks

Our lift of the ADHMN construction is very natural and rather straightforward.

- The **lift of the Dirac operator** was natural considering the corresponding brane configurations.
- It is natural to go to **loop space** to describe self-dual strings.
- The construction nicely involves the **Basu-Harvey equation**.
- It **reduces nicely** to the ADHMN construction.
- The construction does produce **transgressed self-dual strings**.
- A **regression** can be performed to get original self-dual string.
- It can be trivially rendered **nonabelian**:

$$F_{\mu\nu} = \varepsilon_{\mu\nu\rho\sigma} \dot{x}^\rho \nabla^\sigma X^6$$

- Our construction algorithm **generalizes** to the nonabelian case.

The non-abelian tensor multiplet

A recently proposed 3-Lie algebra valued tensor-multiplet implies a transgression.

Recall the **transgression map**:

$$(\mathcal{T}\omega)_x(v_1(\tau), \dots, v_k(\tau)) := \int_{S^1} d\tau \omega(v_1(\tau), \dots, v_k(\tau), \dot{x}(\tau))$$

Equations found by **Lambert, Papageorgakis, 1007.2982**:

$$\nabla^2 X^I - \frac{i}{2}[\bar{\Psi}, \Gamma_\nu \Gamma^I \Psi, C^\nu] - [X^J, C^\nu, [X^J, C_\nu, X^I]] = 0$$

$$\Gamma^\mu \nabla_\mu \Psi - [X^I, C^\nu, \Gamma_\nu \Gamma^I \Psi] = 0$$

$$\nabla_{[\mu} H_{\nu\lambda\rho]} + \frac{1}{4}\varepsilon_{\mu\nu\lambda\rho\sigma\tau}[X^I, \nabla^\tau X^I, C^\sigma] + \frac{i}{8}\varepsilon_{\mu\nu\lambda\rho\sigma\tau}[\bar{\Psi}, \Gamma^\tau \Psi, C^\sigma] = 0$$

$$F_{\mu\nu} - D(C^\lambda, H_{\mu\nu\lambda}) = 0$$

$$\nabla_\mu C^\nu = D(C^\mu, C^\nu) = 0$$

$$D(C^\rho, \nabla_\rho X^I) = D(C^\rho, \nabla_\rho \Psi) = D(C^\rho, \nabla_\rho H_{\mu\nu\lambda}) = 0$$

Factorization of $C^\rho = C\dot{x}^\rho$. Here, **3-Lie algebra transgression**:

$$(\mathcal{T}\omega)_x(v_1(\tau), \dots, v_k(\tau)) := \int_{S^1} d\tau D(\omega(v_1(\tau), \dots, v_k(\tau), \dot{x}(\tau)), C)$$

The non-abelian tensor multiplet on loop space

The corresponding equations can all be rewritten on loop space.

Transgression of matter fields (not in [Huang, Huang, 1008.3834](#))

$$\dot{X}^I(x(\tau)) := R D(C, X^I(x(\tau))), \quad \dot{\Psi}(x(\tau)) := \Gamma^\rho \dot{x}_\rho D(C, \Psi(x(\tau)))$$

Equations of motion (SYM-like):

$$\begin{aligned} \nabla^2 \dot{X}^I + \frac{i}{2} \frac{1}{R} \dot{x}^\nu [\bar{\dot{\Psi}}, \Gamma_\nu \Gamma^I \dot{\Psi}] + [\dot{X}^J, [\dot{X}^J, \dot{X}^I]] &= 0, \\ \frac{1}{R} \Gamma^{\mu\nu} \dot{x}_\nu \nabla_\mu \dot{\Psi} - \Gamma^I [\dot{X}^I, \dot{\Psi}] &= 0, \\ \nabla_\mu \dot{F}^{\mu\nu} + [\dot{X}^I, \nabla^\nu \dot{X}^I] + \frac{i}{2} \left([\bar{\dot{\Psi}}, \Gamma^\nu \dot{\Psi}] - \frac{2}{R^2} \dot{x}^\sigma \dot{x}^\nu [\bar{\dot{\Psi}}, \Gamma_\sigma \dot{\Psi}] \right) &= 0, \end{aligned}$$

Supersymmetry transformations (SYM-like):

$$\begin{aligned} \delta \dot{X}^I &= \frac{1}{R} i \bar{\varepsilon} \Gamma^I \dot{x}^\rho \Gamma_\rho \dot{\Psi}, \\ \delta \dot{A}_\mu &= \frac{1}{R^2} i \bar{\varepsilon} \Gamma_{\mu\lambda} \Gamma_\rho \dot{x}^\lambda \dot{x}^\rho \dot{\Psi}, \\ \delta \dot{\Psi} &= \frac{1}{R} \Gamma^{\nu\mu} \dot{x}_\nu \Gamma^I \nabla_\mu \dot{X}^I \varepsilon + \frac{1}{2} \Gamma_{\mu\nu} \dot{F}^{\mu\nu} \varepsilon - \frac{1}{2} \Gamma^{IJ} [\dot{X}^I, \dot{X}^J] \varepsilon, \\ \delta C^\mu &= 0 \end{aligned}$$

Remarks

The loop space tensor multiplet fits well into the picture.

- Get **SYM theory on loop space** from the tensor multiplet
- **C**-field blocks modes of the theory, need to get rid of it
- Our loop space self-dual string equation **extends compatibly**:

$$F_{\mu\nu} = \varepsilon_{\mu\nu\rho\sigma} \dot{x}^\rho \nabla^\sigma X^6$$

- ADHMN construction for **two M5-branes** using this equation
- **extension** to real and hermitian 3-algebra (ABJM model)
- **Right direction**, more work in progress to get rid of \dot{x} etc.

Conclusions

Summary and Outlook.

Summary:

- ✓ Reformulation of self-dual string equation on **loop space**
- ✓ **Generalized ADHMN construction** for self-dual string
- ✓ **Explicit constructions** in many cases
- ✓ Reformulate **non-abelian tensor multiplet** eqns. on loop space
- ✓ **Partially** generalized ADHMN construction

Future directions:

- ▷ Find **(2,0)-theory**
- ▷ Study **classical integrability** in more detail
- ▷ Quantization of S^3 via **gerbes** and **groupoids**

Describing M5-branes using Loop Spaces

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